

Mimetic Spectral Element Methods

With Hybridization and Algebraic Dual Basis Functions

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MS1001A: Domain Decomposition and Large-Scale Computation

Conservation in physics

In solid mechanics :

- Equilibrium of **forces**
- Equilibrium of **moments**

In fluid mechanics :

- Conservation of **mass**
- Conservation of **vorticity**
- Conservation of **kinetic energy**
- Conservation of **enstrophy**
- Conservation of **helicity**
-

Conservation laws are important.

Example : FVM in CFD.

A fact : In most discretizations, conservations are only satisfied approximately. They will only be exactly satisfied at the limit of mesh refinement (**infinitely refined mesh**).

discrete + **infinitely refined mesh** = continuous level

Mimetic discretizations : aim to satisfy/mimic structures of physics, like conservation laws, at the discrete level.

Structure-preserving discretizations, compatible discretizations.

Structure of PDE

We distinguish two types of basic relations in PDE : **Topological relations** and **Constitutive relations**.

- **Topological relations** : (de Rham complex)

$$\mathbb{R} \hookrightarrow H^1(\Omega) \xrightarrow{\nabla} H(\text{curl}; \Omega) \xrightarrow{\nabla \times} H(\text{div}; \Omega) \xrightarrow{\nabla \cdot} L^2(\Omega) \longrightarrow 0,$$

- **Constitutive relations** : else ; usually related to the constitutive laws.

Example :

Poisson equation : $-\Delta\varphi = f$ or $-\nabla \cdot \nabla\varphi = f$, where $\varphi \in H^1(\Omega)$.

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$$\begin{cases} \mathbf{v} = \nabla\varphi, & \text{so } \mathbf{v} \in H(\text{curl}; \Omega) \\ \nabla \cdot \mathbf{v} = -f, & \text{so } \mathbf{v} \in H(\text{div}; \Omega) \end{cases}$$

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$$\left\{ \begin{array}{l} v = \nabla\varphi, \\ u = \star v, \\ \nabla \cdot u = -f \end{array} \right. \quad \begin{array}{l} \leftarrow \text{Topological relation} \\ \leftarrow \text{Constitutive relation} \\ \leftarrow \text{Topological relation; (fluid : conservation of mass)} \end{array}$$

$v \in H(\text{curl}; \Omega) \xleftrightarrow{\star} u \in H(\text{div}; \Omega)$

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Another example in Linear Elasticity :

constitutive relation $\sigma = C\varepsilon$; topological relation $\nabla \cdot \sigma = -f$ (equilibrium of forces)

Mimetic spectral element method

With **Mimetic spectral element method (MSEM)**^{1, 2, 3}, we preserve the **topological relations** at the discrete level and allow approximation for *constitutive relation*.

Discrete de Rham complex :

$$\begin{array}{ccccccc} \mathbb{R} & \hookrightarrow & H^1(\Omega) & \xrightarrow{\nabla} & H(\text{curl}; \Omega) & \xrightarrow{\nabla \times} & H(\text{div}; \Omega) & \xrightarrow{\nabla \cdot} & L^2(\Omega) & \longrightarrow & 0, \\ & & & & & & \downarrow & & & & \\ \mathbb{R} & \hookrightarrow & H^1(\Omega_h) & \xrightarrow{\nabla} & H(\text{curl}; \Omega_h) & \xrightarrow{\nabla \times} & H(\text{div}; \Omega_h) & \xrightarrow{\nabla \cdot} & L^2(\Omega_h) & \longrightarrow & 0, \end{array}$$

-
1. Kreeft, J., Palha, A. and Gerritsma, M. Mimetic framework on curvilinear quadrilaterals of arbitrary order. *arXiv preprint*, (2011) arXiv :1111.4304.
 2. Kreeft, J. and Gerritsma, M. Mixed mimetic spectral element method for Stokes flow : A pointwise divergence-free solution. *Journal of Computational Physics*, (2013) 240 : 284-309.
 3. Palha, A., Rebelo, P.P., Hiemstra, R., Kreeft, J. and Gerritsma, M. Physics-compatible discretization techniques on single and dual grids, with application to the Poisson equation of volume forms. *Journal of Computational Physics*, (2014) 257 : 1394-1422.

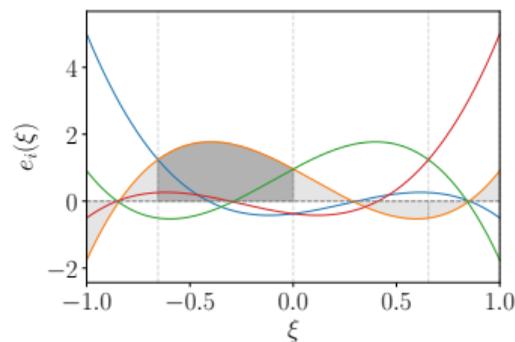
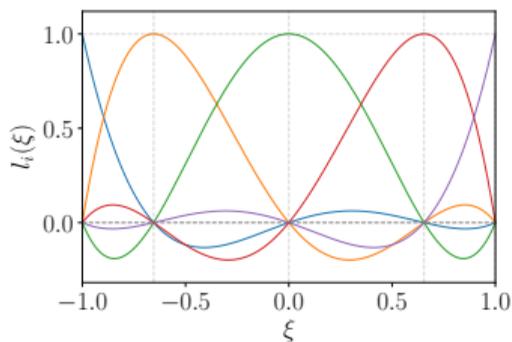
Mimetic spectral element method

Let $-1 \leq \xi_0 < \xi_1 < \dots < \xi_N \leq 1$. The well-known Lagrange polynomials are expressed by $l_i(\xi)$:

$$l_i(\xi) = \prod_{j=0, j \neq i}^N \frac{\xi - \xi_j}{\xi_i - \xi_j}, \quad i \in \{0, 1, 2, \dots, N\}, \text{ satisfying, } l_i(\xi) = \delta_{ij}.$$

The corresponding **edge polynomials**⁴ are

$$e_i(\xi) = - \sum_{k=0}^{i-1} \frac{dl_k(\xi)}{d\xi} = \sum_{k=i}^N \frac{dl_k(\xi)}{d\xi}, \quad i \in \{1, 2, \dots, N\} \text{ satisfying, } \int_{\xi_{j-1}}^{\xi_j} e_i(\xi) = \delta_{ij}.$$



4. Gerritsma, M. Edge functions for spectral element methods. *Spectral and High Order Methods for Partial Differential Equations*. Springer, (2011) 199-207

Mimetic spectral element method

With tensor product, for example in \mathbb{R}^3 ,

$$\mathcal{P} : \{l_i(\xi) \otimes l_j(\eta) \otimes l_k(\zeta)\}$$

$$\mathcal{E} : \{e_i(\xi) \otimes l_j(\eta) \otimes l_k(\zeta), \quad l_i(\xi) \otimes e_j(\eta) \otimes l_k(\zeta), \quad l_i(\xi) \otimes l_j(\eta) \otimes e_k(\zeta)\}$$

$$\mathcal{F} : \{l_i(\xi) \otimes e_j(\eta) \otimes e_k(\zeta), \quad e_i(\xi) \otimes l_j(\eta) \otimes e_k(\zeta), \quad e_i(\xi) \otimes e_j(\eta) \otimes l_k(\zeta)\}$$

$$\mathcal{V} : \{e_i(\xi) \otimes e_j(\eta) \otimes e_k(\zeta)\}$$

$$\text{Discrete de Rham complex :} \quad \mathbb{R} \hookrightarrow \mathcal{P} \xrightarrow{\nabla} \mathcal{E} \xrightarrow{\nabla \times} \mathcal{F} \xrightarrow{\nabla \cdot} \mathcal{V} \longrightarrow 0,$$

While in \mathbb{R}^2 , finite dimensional spaces spanned by basis functions $\{h_i(\xi)e_j(\eta), e_i(\xi)h_j(\eta)\}$ and $\{e_i(\xi)e_j(\eta)\}$ satisfy the **De Rham complex**. Let vector-valued function \mathbf{u} and scalar-valued function f be spanned into

$$\mathbf{u}_h = \left(\sum_{i=0}^N \sum_{j=1}^N u_{i,j} h_i(\xi) e_j(\eta), \sum_{i=1}^N \sum_{j=0}^N v_{i,j} e_i(\xi) h_j(\eta) \right) \quad \text{and} \quad f_h = \sum_{i=1}^N \sum_{j=1}^N f_{i,j} e_i(\xi) e_j(\eta).$$

If $f = \text{div } \mathbf{u}$, then $f_h = \text{div } \mathbf{u}_h$ and

$$f_h = \sum_{i=1}^N \sum_{j=1}^N (u_{i,j} - u_{i-1,j} + v_{i,j} - v_{i,j-1}) e_i(\xi) e_j(\eta),$$

Mimetic spectral element method

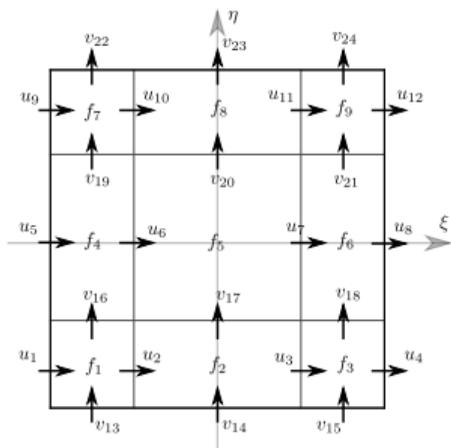


FIGURE – Reference domain.

Let vector-valued function \mathbf{u} and scalar-valued function f be spanned into

$$\mathbf{u}_h = \left(\sum_{i=0}^N \sum_{j=1}^N u_{i,j} h_i(\xi) e_j(\eta), \sum_{i=1}^N \sum_{j=0}^N v_{i,j} e_i(\xi) h_j(\eta) \right), \quad f_h = \sum_{i=1}^N \sum_{j=1}^N f_{i,j} e_i(\xi) e_j(\eta).$$

If $f = \operatorname{div} \mathbf{u}$, then

$$f_h = \operatorname{div} \mathbf{u}_h = \sum_{i=1}^N \sum_{j=1}^N (u_{i,j} - u_{i-1,j} + v_{i,j} - v_{i,j-1}) e_i(\xi) e_j(\eta).$$

Collect all equations and write them in vector form, we have

$$\underline{f} = \mathbb{E}^{2,1} \underline{\mathbf{u}},$$

where

$$\mathbb{E}^{2,1} = \begin{pmatrix} -1 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & -1 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & -1 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & -1 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & -1 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 & 1 & 0 \end{pmatrix}.$$

Mimetic spectral element method

Some literatures :

Kreeft, J. and Gerritsma, M., Mixed mimetic spectral element method for **Stokes flow** : A pointwise divergence-free solution. *Journal of Computational Physics*, (2013) 240 : 284-309.

Palha, A., Rebelo, P.P., Hiemstra, R., Kreeft, J. and Gerritsma, M., Physics-compatible discretization techniques on single and dual grids, with application to the **Poisson equation** of volume forms. *Journal of Computational Physics*, (2014) 257 : 1394-1422.

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Lee D. and Palha A., A mixed mimetic spectral element model of the 3D **compressible Euler equations** on the cubed sphere, *Journal of Computational Physics*, (2019)

(In preparation), A mass, kinetic energy and helicity conserving mimetic spectral element method for **3D incompressible Navier-Stokes equations**.

Hybridization & dual basis functions

Mimetic spectral element method is computationally expensive.

- large systems
 - mixed formulation
 - high order method
- low sparsity

Not ready for large scale computation!

Hybridization & dual basis functions

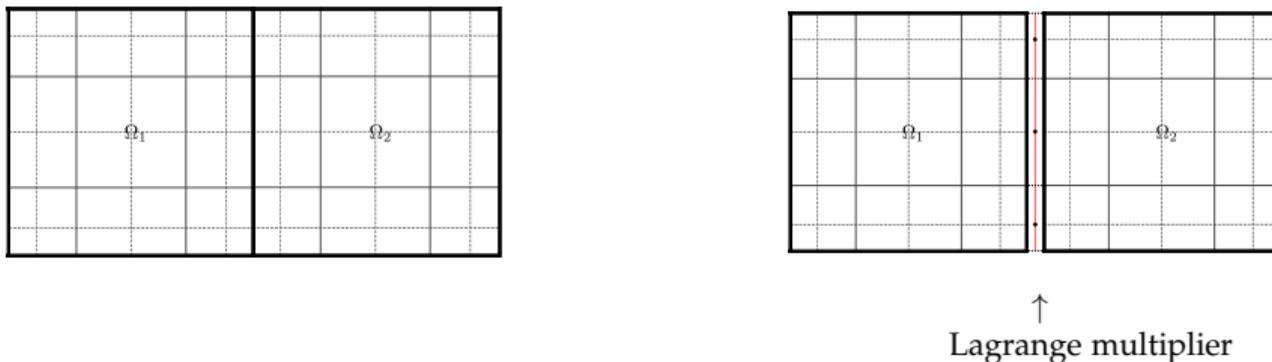
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- large systems ← Domain decomposition : Hybridization
 - mixed formulation
 - high order method
- low sparsity ← Dual basis functions

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Hybridization

Hybrid finite element methods are those methods that first allow the discontinuity across the inter-element interface then re-enforce (weakly or strongly) the continuity by introducing a Lagrange multiplier between elements.



Similar idea has also been used in, e.g., *mortar methods* and *finite element tearing and interconnecting (FETI) methods*, and more.

Dual basis functions

Using dual basis functions eliminates some metric-dependent matrices from the discrete system.

$$(\phi, \varphi) = \underline{\phi}^T \mathbb{M} \underline{\varphi}$$

$$(\phi, \tilde{\varphi}) = \underline{\phi}^T \tilde{\underline{\varphi}}$$

where $\tilde{\varphi}$ represent it is expanded with dual basis functions.

Example : Poisson equation⁵

The well-known mixed formulation of the Poisson equation : Find $(\mathbf{u}, \varphi) \in H(\text{div}, \Omega) \times H^1(\Omega)$ such that

$$\begin{cases} (\mathbf{u}, \mathbf{v}) + (\varphi, \text{div } \mathbf{v}) & = \langle \hat{\varphi}, \text{tr}_{\text{div } \mathbf{v}} \rangle_{\partial\Omega} \\ (\psi, \text{div } \mathbf{u}) & = -(\psi, f) \end{cases},$$

for all $(\mathbf{v}, \psi) \in H(\text{div}, \Omega) \times H^1(\Omega)$. This is a weak mixed formulation of the Poisson equation.

Discretization :

$$\boxed{\begin{pmatrix} \mathbb{M}_1 & (\mathbb{M}_2 \mathbb{E})^T \\ \mathbb{M}_2 \mathbb{E} & 0 \end{pmatrix} \begin{pmatrix} \underline{\mathbf{u}} \\ \underline{\varphi} \end{pmatrix} = \begin{pmatrix} \mathbb{B} \\ -\mathbb{M}_2 \underline{f} \end{pmatrix}.}$$

5. Y. Zhang, V. Jain, A. Palha, M. Gerritsma, The discrete Steklov-Poincare operator using algebraic dual polynomials, Computational Methods in Applied Mathematics, (2019)

Example : Poisson equation, hybridization

If we set up a mesh Ω_h in Ω , by breaking \mathbf{u} and φ into **broken spaces**, $H(\text{div}, \Omega_h)$ and $H^1(\Omega_h)$, and introducing a new Lagrange multiplier $\check{\varphi}$ in the **interface space** $H^{1/2}(\partial\Omega_h \setminus \partial\Omega)$

Given $f \in L^2(\Omega)$ and $\hat{\varphi} = \text{tr}_{\text{grad}} \varphi \in H^{1/2}(\partial\Omega)$, find $(\mathbf{u}, \varphi, \check{\varphi}) \in H(\text{div}, \Omega_h) \times H^1(\Omega_h) \times H^{1/2}(\partial\Omega_h \setminus \partial\Omega)$ such that

$$\begin{cases} (\mathbf{u}, \mathbf{v}) + (\varphi, \text{div } \mathbf{v}) - \langle \check{\varphi}, \text{tr}_{\text{div}} \mathbf{v} \rangle_{\partial\Omega_h \setminus \partial\Omega} & = \langle \hat{\varphi}, \text{tr}_{\text{div}} \mathbf{v} \rangle \\ (\psi, \text{div } \mathbf{u}) & = -(\psi, f) \\ -\langle \check{\psi}, \text{tr}_{\text{div}} \mathbf{u} \rangle_{\partial\Omega_h \setminus \partial\Omega} & = 0 \end{cases},$$

for all $(\mathbf{v}, \psi, \check{\psi}) \in H(\text{div}, \Omega_h) \times H^1(\Omega_h) \times H^{1/2}(\partial\Omega_h \setminus \partial\Omega)$.

Discretization :

$$\begin{pmatrix} \mathbf{M}_1 & \mathbf{M}_2 \mathbb{E}^T & -(\mathbf{M} \mathbf{N}_I)^T \\ \mathbf{M}_2 \mathbb{E} & 0 & 0 \\ -\mathbf{M} \mathbf{N}_I & 0 & 0 \end{pmatrix} \begin{pmatrix} \underline{\mathbf{u}} \\ \underline{\varphi} \\ \underline{\check{\varphi}} \end{pmatrix} = \begin{pmatrix} \mathbb{B} \hat{\varphi} \\ -\mathbf{M}_2 f \\ 0 \end{pmatrix}.$$

Example : Poisson equation, hybridization + dual basis functions

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Discretization :

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Example : Poisson equation, hybridization + dual basis functions

Discrete hybrid mixed formulation :

$$\begin{pmatrix} \mathbf{M}_1 & \mathbf{E}^T & -\mathbf{N}_I^T \\ \mathbf{E} & 0 & 0 \\ -\mathbf{N}_I & 0 & 0 \end{pmatrix} \begin{pmatrix} \underline{\mathbf{u}} \\ \underline{\varphi} \\ \underline{\tilde{\varphi}} \end{pmatrix} = \begin{pmatrix} \mathbf{N}_B^T \hat{\varphi} \\ -\underline{f} \\ 0 \end{pmatrix}.$$

- \mathbf{M}_1 : metric-dependent ; element-wise-block-diagonal ;
- \mathbf{E} : metric-independent ; element-wise-block-diagonal ; super sparse ; ± 1 non-zero entries ;
- \mathbf{N} : metric-independent ; even more sparse ; ± 1 non-zero entries ;

$$\mathbf{E} = \begin{pmatrix} -1 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & -1 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & -1 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & -1 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & -1 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & -1 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 & 1 \end{pmatrix}.$$

Example : Poisson equation, hybridization + dual basis functions

Discrete hybrid mixed formulation :

$$\begin{pmatrix} \mathbf{M}_1 & \mathbf{E}^T & -\mathbf{N}_I^T \\ \mathbf{E} & 0 & 0 \\ -\mathbf{N}_I & 0 & 0 \end{pmatrix} \begin{pmatrix} \underline{\mathbf{u}} \\ \underline{\varphi} \\ \underline{\check{\varphi}} \end{pmatrix} = \begin{pmatrix} \mathbf{N}_B^T \hat{\varphi} \\ -\underline{f} \\ 0 \end{pmatrix}.$$

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- \mathbf{E} : metric-independent ; element-wise-block-diagonal ; super sparse ; ± 1 non-zero entries ;
- \mathbf{N} : metric-independent ; even more sparse ; ± 1 non-zero entries ;

A reduced system for the interface variable $\underline{\check{\varphi}}$:

$$\mathbf{S} \underline{\check{\varphi}} = \mathbf{F},$$

where

$$\mathbf{S} = -\mathbf{N}_I \mathbf{M}_1^{-1} \left[\mathbf{M}_1 - \mathbf{E}^T \left(\mathbf{E} \mathbf{M}_1^{-1} \mathbf{E}^T \right)^{-1} \mathbf{E} \right] \mathbf{M}_1^{-1} \mathbf{N}_I^T,$$

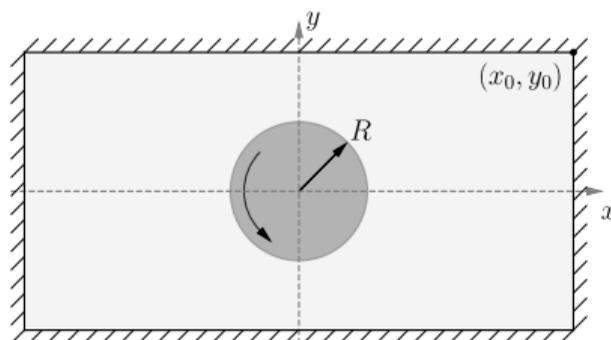
$$\mathbf{F} = \mathbf{F}_{\hat{\varphi}} + \mathbf{F}_f,$$

$$\mathbf{F}_{\hat{\varphi}} = \mathbf{N}_I \mathbf{M}_1^{-1} \left[\mathbf{M}_1 - \mathbf{E}^T \left(\mathbf{E} \mathbf{M}_1^{-1} \mathbf{E}^T \right)^{-1} \mathbf{E} \right] \mathbf{M}_1^{-1} \mathbf{N}_B^T \hat{\varphi},$$

$$\mathbf{F}_f = -\mathbf{N}_I \mathbf{M}_1^{-1} \mathbf{E}^T \left(\mathbf{E} \mathbf{M}_1^{-1} \mathbf{E}^T \right)^{-1} \underline{f}.$$

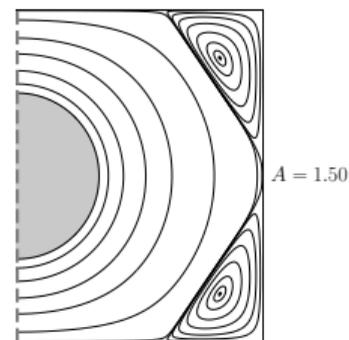
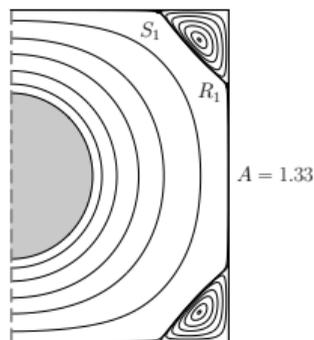
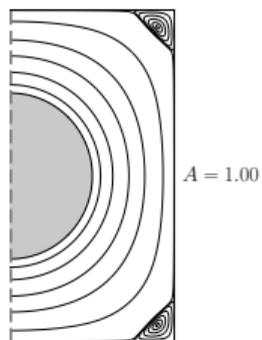
- Inverting $\mathbf{M}^{(1)}$ and $\mathbf{E}^{2,1} \mathbf{M}^{(1)^{-1}} \mathbf{E}^{2,1T}$ is easy (in parallel) because they are element-wise-block-diagonal.
- Solving for $\underline{\check{\varphi}}$ is cheap (smaller system size and condition number). Remaining local problems for $\underline{\mathbf{u}}$ and $\underline{\varphi}$ are trivial.

Stokes flow induced by cylinder rotation

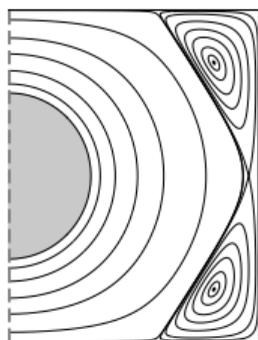
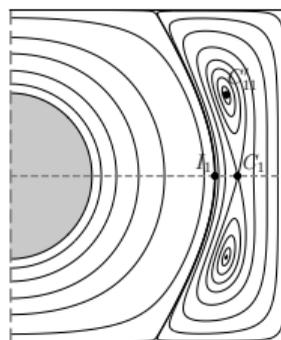
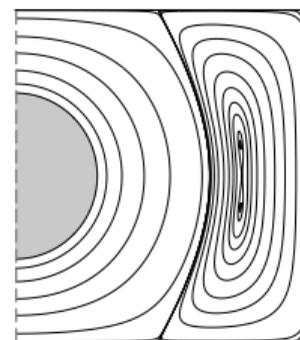


$$A = \frac{x_0}{y_0}$$

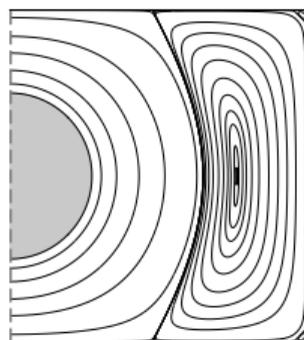
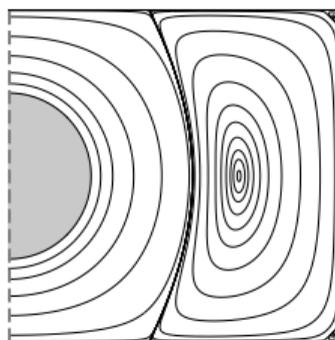
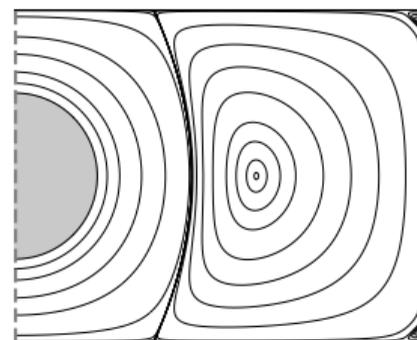
Stokes flow induced by cylinder rotation



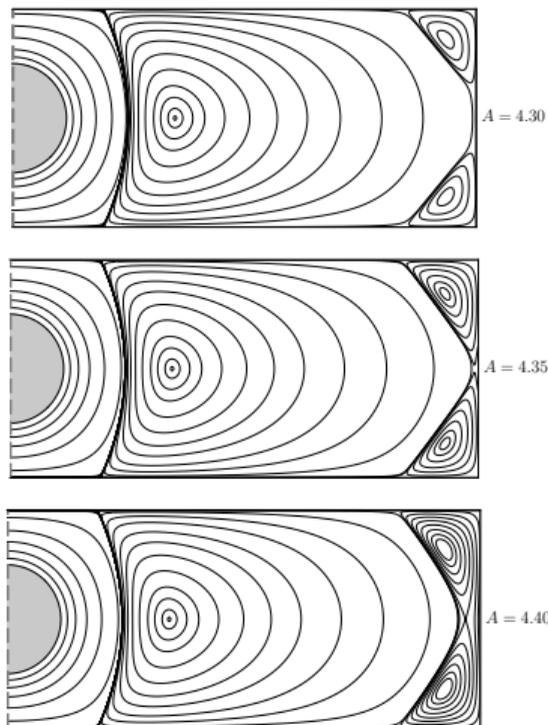
Stokes flow induced by cylinder rotation


 $A = 1.53$

 $A = 1.68$

 $A = 1.79$

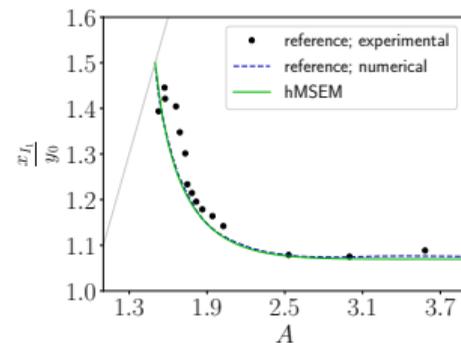
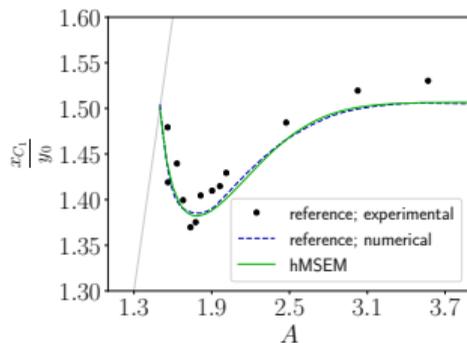
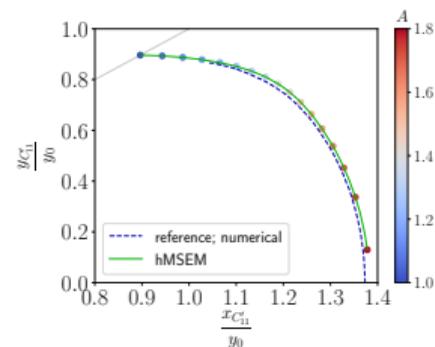
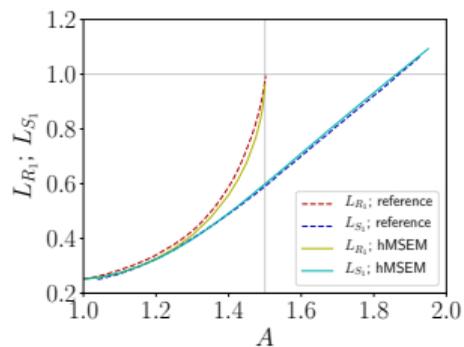
Stokes flow induced by cylinder rotation

 $A = 1.81$  $A = 2.02$  $A = 2.51$

Stokes flow induced by cylinder rotation



Stokes flow induced by cylinder rotation



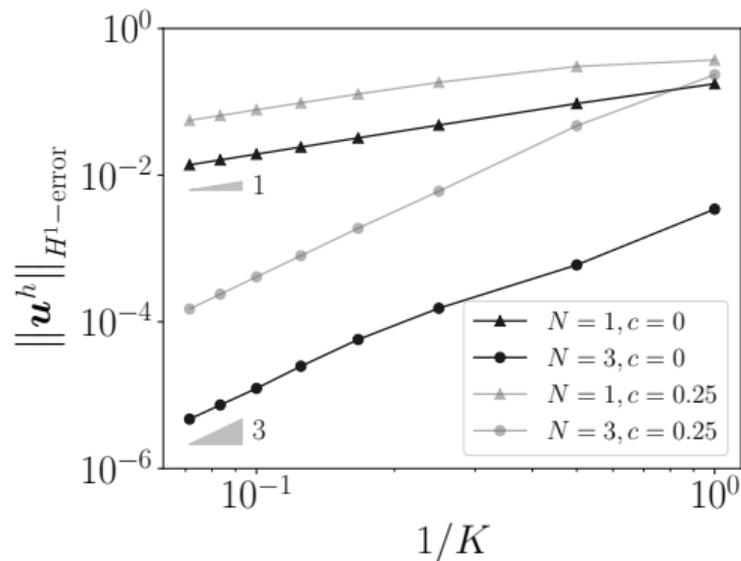
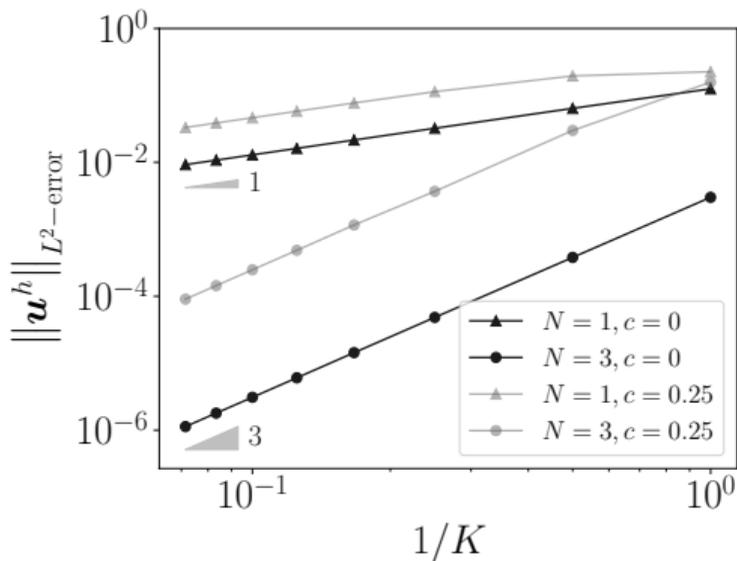
Accuracy

TABLE – Results of $\|x\|_{L^2\text{-error}}$ and $\|x - x'\|_{L^2\text{-norm}}$ (in brackets), where x and x' are solutions of the hMSEM and MSEM respectively, for $N \in \{1, 3\}$, $K \in \{2, 4, 6\}$, and $c \in \{0, 0.25\}$.

x	K	$N = 1$		$N = 3$	
		$c = 0$	$c = 0.25$	$c = 0$	$c = 0.25$
u^h	2	6.4029E-2(2.05E-16)	1.9534E-1(2.79E-16)	3.8024E-4(5.57E-16)	2.9940E-2(1.70E-15)
	4	3.2265E-2(5.40E-16)	1.1353E-1(4.61E-16)	4.8312E-5(5.67E-15)	3.6850E-3(1.48E-14)
	6	2.1542E-2(7.03E-16)	7.7069E-2(8.71E-16)	1.4377E-5(4.08E-14)	1.1604E-3(1.06E-13)
ω^h	2	4.7436E-2(6.76E-16)	5.1309E-2(9.59E-16)	2.8846E-4(1.22E-14)	1.2456E-2(3.17E-14)
	4	2.3986E-2(2.98E-15)	3.8150E-2(3.31E-15)	8.2990E-5(2.63E-13)	1.0417E-3(1.00E-12)
	6	1.6008E-2(2.10E-14)	2.2952E-2(2.25E-14)	3.0156E-5(4.08E-12)	2.3494E-4(1.02E-11)
σ^h	2	9.3659E-1(1.84E-14)	2.6588(3.13E-14)	5.6919E-3(2.21E-13)	4.5920E-1(4.12E-13)
	4	4.5869E-1(8.16E-14)	1.6633(6.25E-14)	1.6879E-3(4.03E-12)	6.1945E-2(1.48E-11)
	6	3.0391E-1(2.58E-13)	1.1638(2.82E-13)	6.2626E-4(5.39E-11)	1.9624E-2(1.43E-10)

Accuracy

Super-convergence with respect to H^1 -error :



Conclusions

We present the **mimetic spectral element method** and its hybridization with dual basis functions.

- It is **mimetic**/structure-preserving; it uses nodal and integral values (with respect to nodes, edges, faces and volumes) as dof's.
- It is a **hybrid**/domain decomposition method;
- It uses **dual polynomials**; some discrete matrices are **metric-free**, **extremely sparse** and **low order finite-difference(volume)-like** (containing non-zero entries of -1 and 1 only).

Thanks a lot. Questions?

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