




Discrete Geometries of Mathematics and Physics

# Mimetic spectral element method<sup>1</sup>

## Assignment #2

From the reference element to physical elements

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We have known how to set up mimetic spectral element spaces in the reference element  $\Omega_r = [-1, 1]^2$  through Assignment #1. However, for a real problem, its computation domain most likely is different to  $\Omega_r$ . Thus we need to know how to construct mimetic spectral element spaces in a more general element such that we can apply the mimetic spectral element method to a real problem.

### 1 Coordinate transformation

Consider a domain  $\Omega_0$  in the coordinate system  $(\xi, \eta)$  and a domain  $\Omega$  in the coordinate system  $(x, y)$ . A  $C^1$  diffeomorphism (both itself and its inverse are  $C^1$  continuous)  $\Phi$  maps  $\Omega_0$  into  $\Omega$ , i.e.,

$$\Phi : \Omega_0 \rightarrow \Omega,$$

by

$$\begin{bmatrix} x \\ y \end{bmatrix} = \Phi(\xi, \eta) = \begin{bmatrix} \Phi_x(\xi, \eta) \\ \Phi_y(\xi, \eta) \end{bmatrix}$$

Since the mapping  $\Phi$  is  $C^1$  continuous, we can compute its Jacobian matrix  $\mathcal{J}$ ,

$$\mathcal{J} = \begin{bmatrix} \mathcal{J}_{11} & \mathcal{J}_{12} \\ \mathcal{J}_{21} & \mathcal{J}_{22} \end{bmatrix} = \begin{bmatrix} \frac{\partial x}{\partial \xi} & \frac{\partial x}{\partial \eta} \\ \frac{\partial y}{\partial \xi} & \frac{\partial y}{\partial \eta} \end{bmatrix}.$$

The Jacobian,  $J$ , of the mapping is the determinant of the Jacobian matrix, i.e.,

$$J = \det(\mathcal{J}).$$

<sup>1</sup>[https://mathischeap.com/contents/teaching/advanced\\_numerical\\_methods/mimetic\\_spectral\\_element\\_method/main](https://mathischeap.com/contents/teaching/advanced_numerical_methods/mimetic_spectral_element_method/main)

The metric matrix is defined as

$$\mathcal{G} = \begin{bmatrix} g_{11} & g_{12} \\ g_{21} & g_{22} \end{bmatrix} := \mathcal{J}^\top \mathcal{J},$$

To be more exact,  $g_{ij} = \sum_{m=1}^2 \mathcal{J}|_{mi} \mathcal{J}|_{mj}$ . For example,  $g_{11} = \mathcal{J}_{11}\mathcal{J}_{11} + \mathcal{J}_{21}\mathcal{J}_{21}$  and  $g_{12} = \mathcal{J}_{11}\mathcal{J}_{12} + \mathcal{J}_{21}\mathcal{J}_{22}$ . The determinate of the metric matrix is called the metric,  $g$ , of the mapping, namely,

$$g := \det(\mathcal{G}).$$

One can show that the metric is equal to the square of the Jacobian, i.e.,

$$g = J^2 = [\det(\mathcal{J})]^2.$$

As  $\Phi$  is a  $C^1$  diffeomorphism, its inverse mapping  $\Phi^{-1}$ ,

$$\Phi^{-1} : \Omega \rightarrow \Omega_0,$$

must also be  $C^1$  continuous. The Jacobian matrix of  $\Phi^{-1}$  is called the inverse Jacobian matrix of  $\Phi$  is

$$\mathcal{J}^{-1} := \begin{bmatrix} \mathcal{J}_{11}^{-1} & \mathcal{J}_{12}^{-1} \\ \mathcal{J}_{21}^{-1} & \mathcal{J}_{22}^{-1} \end{bmatrix} = \begin{bmatrix} \frac{\partial \xi}{\partial x} & \frac{\partial \xi}{\partial y} \\ \frac{\partial \eta}{\partial x} & \frac{\partial \eta}{\partial y} \end{bmatrix}.$$

It must satisfy  $\mathcal{J}^{-1} \mathcal{J} = \mathcal{I}$  where  $\mathcal{I}$  is the identity matrix. The inverse Jacobian is the determinant of  $\mathcal{J}^{-1}$ , i.e.,

$$J^{-1} = \det(\mathcal{J}^{-1}) = \frac{1}{J} = \frac{1}{\det(\mathcal{J})}.$$

The inverse metric matrix is defined as

$$\mathcal{G}^{-1} = \begin{bmatrix} g^{11} & g^{12} \\ g^{21} & g^{22} \end{bmatrix} := \mathcal{J}^{-1} (\mathcal{J}^{-1})^\top,$$

where  $g^{ij} = \sum_{m=1}^2 \mathcal{J}^{-1}|_{im} \mathcal{J}^{-1}|_{jm}$ . For example,  $g^{11} = \mathcal{J}_{11}^{-1}\mathcal{J}_{11}^{-1} + \mathcal{J}_{12}^{-1}\mathcal{J}_{12}^{-1}$  and  $g^{12} = \mathcal{J}_{11}^{-1}\mathcal{J}_{21}^{-1} + \mathcal{J}_{12}^{-1}\mathcal{J}_{22}^{-1}$ .

## 2 Transformation of mimetic spectral element spaces

We now consider a  $C^1$  diffeomorphism mapping that maps the reference element into a particular element  $\Omega_n$ , called the physical element, i.e.,

$$\Phi_n : \Omega_r \rightarrow \Omega_n.$$

What we really care about is how the mimetic spectral element spaces are transformed under this mapping. Once we know it, we can set up mimetic spectral element spaces in a physical element rather than in the reference element only.

In this assignment, to limit the difficulty at a reasonable level, we will only consider the physical element to be an orthogonal (whose edges are parallel to axes) rectangle, i.e.,

$$\Omega_n := [x_0, x_1] \times [y_0, y_1],$$

where  $x_0, x_1, y_0, y_1 \in \mathbb{R}$  and  $x_0 < x_1, y_0 < y_1$ . With such a physical, we know the mapping can be explicitly expressed as

$$\begin{bmatrix} x \\ y \end{bmatrix} = \Phi_n(\xi, \eta) = \begin{bmatrix} \Phi_{n,x}(\xi) \\ \Phi_{n,y}(\eta) \end{bmatrix} = \begin{bmatrix} x_0 + \frac{d_x(\xi+1)}{2} \\ y_0 + \frac{d_y(\eta+1)}{2} \end{bmatrix},$$

where  $d_x = x_1 - x_0$  and  $d_y = y_1 - y_0$ . And we can easily compute its metric-related values,

$$\begin{aligned}\mathcal{J}_n &= \begin{bmatrix} \frac{d_x}{2} & 0 \\ 0 & \frac{d_y}{2} \end{bmatrix}, \\ J_n &= \det(\mathcal{J}_n) = \frac{d_x d_y}{4}, \\ \mathcal{G}_n &= \begin{bmatrix} \frac{d_x^2}{4} & 0 \\ 0 & \frac{d_y^2}{4} \end{bmatrix}, \\ g_n &= \det(\mathcal{G}_n) = J_n^2 = \frac{d_x^2 d_y^2}{16}, \\ \mathcal{J}_n^{-1} &= \begin{bmatrix} \frac{2}{d_x} & 0 \\ 0 & \frac{2}{d_y} \end{bmatrix}, \\ J_n^{-1} &= \frac{4}{d_x d_y}, \\ \mathcal{G}_n^{-1} &= \begin{bmatrix} \frac{4}{d_x^2} & 0 \\ 0 & \frac{4}{d_y^2} \end{bmatrix}.\end{aligned}$$

Note that these metric-related values are all constants and are independent to the origin of the rectangle, i.e.  $(x_0, y_0)$ . That means if rectangles are of the same shape, they will all have same metric-related values.

To obtain transformed mimetic spectral element spaces in  $\Omega_n$ , we just need to find the transformed basis functions. And if these basis functions are still linearly independent, they span the mimetic spectral element spaces in  $\Omega_n$ . So, after this point, we introduce how to transform the basis functions.

## 2.1 Transformation: $C(\Omega_r) \rightarrow C(\Omega_n)$ (or $G(\Omega_r) \rightarrow G(\Omega_n)$ )

We have learned from Assignment #1 that the basis functions of mimetic spectral element space  $C(\Omega_r)$  are

$$\mathbf{11}^{ij}(\xi, \eta) = l^i(\xi)l^j(\eta), \quad i, j \in \{0, 1, \dots, N\}.$$

We can transform them to  $\Omega_n$  as

$$\mathbf{11}_n^{ij}(x, y) = \mathbf{11}_n^{ij}(\Phi(\xi, \eta)) = \mathbf{11}^{ij}(\xi, \eta), \quad i, j \in \{0, 1, \dots, N\}.$$

They are a set of basis functions for the mimetic spectral element space  $C(\Omega_n)$  and  $G(\Omega_n)$ . So, if  $\omega_h \in C(\Omega_n)$  or  $G(\Omega_n)$ , we can express it as

$$(1) \quad \omega_h = \sum_{i=0}^N \sum_{j=0}^N \mathbf{w}_{ij} \mathbf{11}_n^{ij}(x, y).$$

Given a function  $\omega(x, y)$  in  $\Omega_n$ , if  $\omega_h = \pi(\omega)$ , the expansion coefficients are

$$(2) \quad \mathbf{w}_{ij} = \omega(x_i, y_j) = \omega(\Phi_n(\xi_i, \eta_j)).$$

## 2.2 Transformation: $\mathbf{D}(\Omega_r) \rightarrow \mathbf{D}(\Omega_n)$

We can transform basis functions of  $\mathbf{D}(\Omega_r)$  to  $\Omega_n$  by

$$\begin{aligned} \mathbf{le}_n^{ij}(x, y) &= \frac{\mathcal{J}_n|_{11}}{J_n} \mathbf{le}^{ij}(\Phi_n^{-1}(x, y)) = \frac{2}{d_y} \mathbf{le}^{ij}(\Phi_n^{-1}(x, y)) = \frac{2}{d_y} \mathbf{le}^{ij}(\xi, \eta), \\ \mathbf{el}_n^{ij}(x, y) &= \frac{\mathcal{J}_n|_{22}}{J_n} \mathbf{el}^{ij}(\Phi_n^{-1}(x, y)) = \frac{2}{d_x} \mathbf{el}^{ij}(\Phi_n^{-1}(x, y)) = \frac{2}{d_x} \mathbf{el}^{ij}(\xi, \eta), \end{aligned}$$

They are basis functions that span the mimetic spectral element space  $\mathbf{D}(\Omega_n)$ . If  $\mathbf{u}_h \in \mathbf{D}(\Omega_n)$ , we can express it as

$$(3) \quad \mathbf{u}_h = \begin{bmatrix} \sum_{i=0}^N \sum_{j=1}^N \mathbf{u}_{ij} \mathbf{le}_n^{ij}(x, y) \\ \sum_{i=1}^N \sum_{j=0}^N \mathbf{v}_{ij} \mathbf{el}_n^{ij}(x, y) \end{bmatrix}$$

Given a vector  $\mathbf{u}(x, y) = \begin{bmatrix} u(x, y) \\ v(x, y) \end{bmatrix}$  in  $\Omega_n$ , if  $\mathbf{u}_h = \pi(\mathbf{u})$ , the expansion coefficients of  $\mathbf{u}_h$  are

$$(4) \quad \begin{aligned} \mathbf{u}_{ij} &= \int_{y_{j-1}}^{y_j} u(x_i, y) dy = \int_{\Phi_{n,y}(\eta_{j-1})}^{\Phi_{n,y}(\eta_j)} u(\Phi_{n,x}(\xi_i), y) dy \\ \mathbf{v}_{ij} &= \int_{x_{i-1}}^{x_i} v(x, y_j) dx = \int_{\Phi_{n,x}(\xi_{i-1})}^{\Phi_{n,x}(\xi_i)} v(x, \Phi_{n,y}(\eta_j)) dx \end{aligned}$$

## 2.3 Transformation: $\mathbf{R}(\Omega_r) \rightarrow \mathbf{R}(\Omega_n)$

In orthogonal rectangle elements, the basis functions  $\mathbf{R}(\Omega_n)$  are same to those of  $\mathbf{D}(\Omega_n)$ . But they span different components. In details, if  $\boldsymbol{\sigma}_h \in \mathbf{R}(\Omega_n)$ , we can express it as

$$(5) \quad \boldsymbol{\sigma}_h = \begin{bmatrix} \sum_{i=1}^N \sum_{j=0}^N \mathbf{s}_{ij} \mathbf{el}_n^{ij}(x, y) \\ \sum_{i=0}^N \sum_{j=1}^N \mathbf{t}_{ij} \mathbf{le}_n^{ij}(x, y) \end{bmatrix}$$

And given a vector  $\boldsymbol{\sigma}(x, y) = \begin{bmatrix} s(x, y) \\ t(x, y) \end{bmatrix}$  in  $\Omega_n$ , if  $\boldsymbol{\sigma}_h = \pi(\boldsymbol{\sigma})$ , the expansion coefficients of  $\boldsymbol{\sigma}_h$  are

$$(6) \quad \begin{aligned} \mathbf{s}_{ij} &= \int_{x_{i-1}}^{x_i} s(x, y_j) dx = \int_{\Phi_{n,x}(\xi_{i-1})}^{\Phi_{n,x}(\xi_i)} s(x, \Phi_{n,y}(\eta_j)) dx \\ \mathbf{t}_{ij} &= \int_{y_{j-1}}^{y_j} t(x_i, y) dy = \int_{\Phi_{n,y}(\eta_{j-1})}^{\Phi_{n,y}(\eta_j)} t(\Phi_{n,x}(\xi_i), y) dy \end{aligned}$$

## 2.4 Transformation: $S(\Omega_r) \rightarrow S(\Omega_n)$

Fianlly, we know from Assigment #1 that the basis functions of mimetic spectral element space  $S(\Omega_r)$  are

$$\mathbf{ee}^{ij}(\xi, \eta) = e^i(\xi) e^j(\eta), \quad i, j \in \{1, 2, \dots, N\}.$$

We can transform them to  $\Omega_n$  as

$$\mathbf{ee}_n^{ij}(x, y) = \frac{1}{J_n} \mathbf{ee}^{ij}(\Phi_n^{-1}(x, y)) = \frac{4}{d_x d_y} \mathbf{ee}^{ij}(\xi, \eta)$$

They are a set of basis functions for the mimetic spectral element space  $S(\Omega_n)$ . Thus, if  $\phi_h \in S(\Omega_n)$ , we can express it as

$$(7) \quad \phi_h = \sum_{i=1}^N \sum_{j=1}^N f_{ij} \mathbf{e}_n^{ij}(x, y).$$

Given a scalar  $\phi(x, y)$  in  $\Omega_n$ , if  $\phi_h = \pi(\phi)$ , the expansion coefficients of  $\phi_h$  are

$$(8) \quad f_{ij} = \int_{x_{i-1}}^{x_i} \int_{y_{j-1}}^{y_j} \phi(x, y) dy dx = \int_{\Phi_{n,x}(\xi_{i-1})}^{\Phi_{n,x}(\xi_i)} \int_{\Phi_{n,y}(\eta_{j-1})}^{\Phi_{n,y}(\eta_j)} \phi(x, y) dy dx.$$

#### Assignment 2.1.0: Reduction and reconstruction of $C(\Omega_n)$ and $G(\Omega_n)$

You need to program two functions to compute the reduction (2) and the reconstruction (1) for  $C(\Omega_n)$  and  $G(\Omega_n)$ , respectively.

#### Assignment 2.1.1: Reduction and reconstruction of $D(\Omega_n)$

You need to program two functions to compute the reduction (4) and the reconstruction (3) for  $D(\Omega_n)$ , respectively.

#### Assignment 2.1.2: Reduction and reconstruction of $R(\Omega_n)$

You need to program two functions to compute the reduction (6) and the reconstruction (5) for  $R(\Omega_n)$ , respectively.

#### Assignment 2.1.3: Reduction and reconstruction of $S(\Omega_n)$

You need to program two functions to compute the reduction (8) and the reconstruction (7) for  $S(\Omega_n)$ , respectively.

#### Assignment 2.1.4: Incidence matrices

You now need to use the incidence matrices you have programmed in Assignment #1 to this assignment and check out if it still works.

The answer should be "YES". This means we no need to compute new incidence matrices for the physical domain  $\Omega_n$ . Think about it. Why?

### 3 Three dimensions

Read [1, Chapter 2] for the more details of three-dimensional cases. They can help you understanding two-dimensional cases. Here in this assignment, we only present the necessary contents for

programming the functionalities. Reading more insights can help you understanding the idea and philosophy behind very much.

## References

- [1] Y. Zhang, Mimetic Spectral Element Method and Extensions toward Higher Computational Efficiency (2022).